

# The equation of state problem

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## effective theories and first principles

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Department of Physics  
University of Surrey

UK LOFT meeting, 24 June 2013



# Equation of State

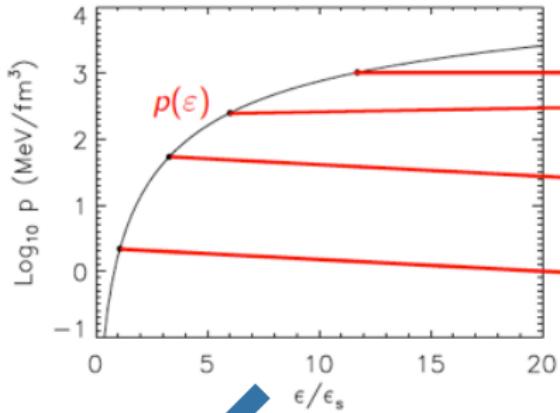
Tolman-Oppenheimer-Volkov equations

$$\frac{dP}{dr} = -\frac{G}{c^2} \frac{(m + 4\pi Pr^3)(\epsilon + P)}{r(r - 2Gm/c^2)} \quad \frac{dm}{dr} = \frac{4\pi\epsilon r^2}{c^2}$$



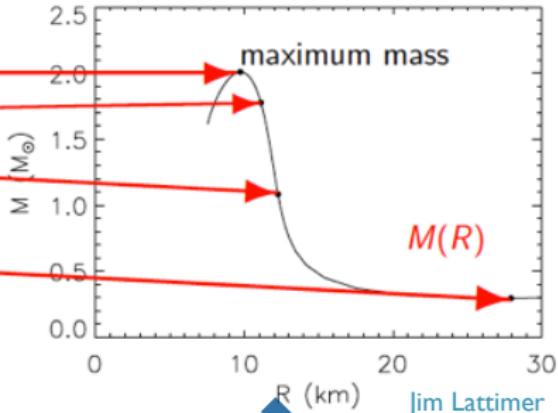
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## Equation of State



Nuclear quantity  
(that's what I do!)

## Mass-Radius relation



Stellar quantity

# Equation of State

Tolman-Oppenheimer-Volkov equations

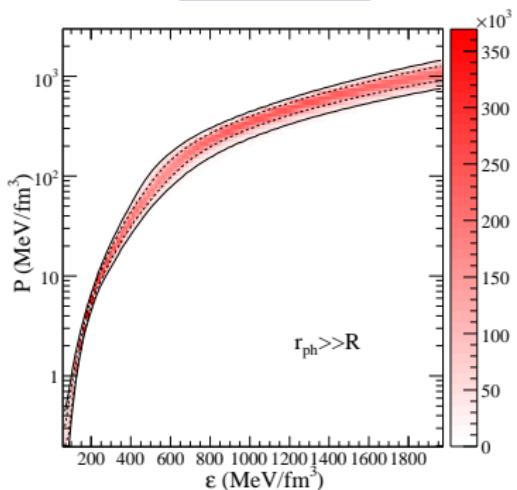
$$\frac{dP}{dr} = -\frac{G}{c^2} \frac{(m + 4\pi Pr^3)(\epsilon + P)}{r(r - 2Gm/c^2)}$$

$$\frac{dm}{dr} = \frac{4\pi\epsilon r^2}{c^2}$$



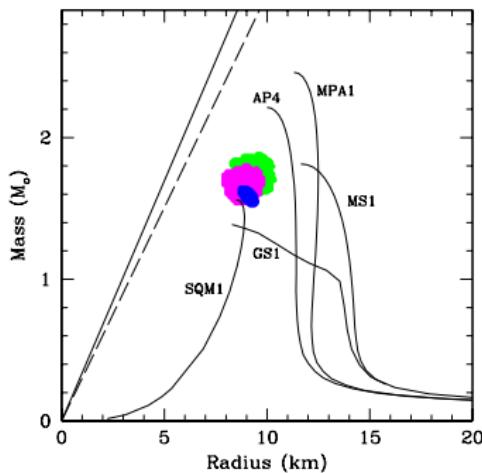
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## Empirical EoS



Steiner *et al.*, ApJ 722, 33 (2010)

## Mass-radius relation



Özel, Baym & Güver, PRD 82, 101301R (2010)

- Mass-Radius relation from bursts
- Bayesian data analysis to get model-independent EoS
  - ① 3 type-I X-ray bursts
  - ② 3 transient low mass X-ray binaries
  - ③ 1 isolated cooling NS, RX J1856-3754



# Isospin asymmetric matter

Tuning correlations

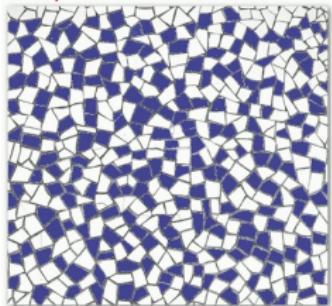


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## Nuclear “trencadís”

$\beta=0$

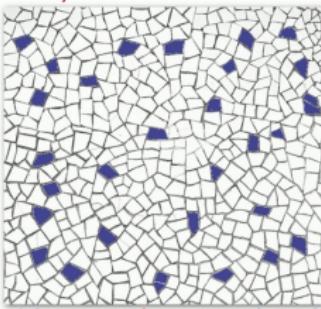
Symmetric matter



SR+Tensor correlations

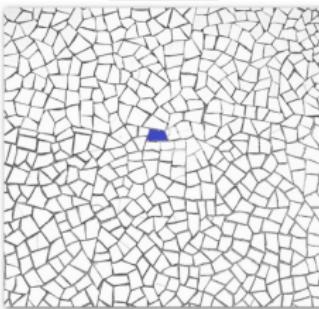
$\beta \neq 0$

Asymmetric matter



Neutrons **less** correlated  
Protons **more** correlated

$\beta \approx 1$   
Polaron



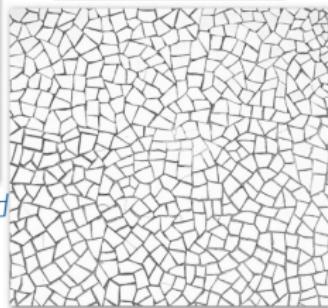
Protons **maximally** correlated  
**Hyper-impurities?**

Neutron stars



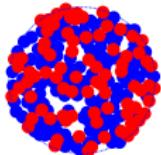
$\beta=1$

Neutron matter



SR correlations

Nuclei



$$\beta = \frac{N - Z}{N + Z}$$

- Frick, Rios et al. PRC **71**, 014313 (2005)  
Rios et al. PRC **79**, 064308 (2009)  
Carbone et al. EPL **97** 22001 (2012)

# Effective approach to the EoS

- EoS provides a characterization of bulk properties:

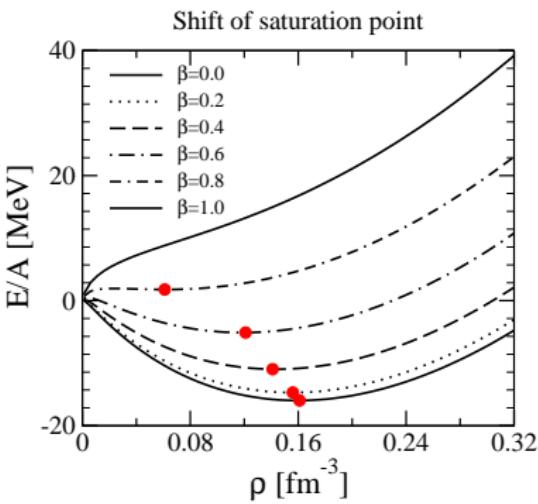
$$p(\varepsilon) =? \iff \frac{E}{A}(\rho, \beta) =?$$

$$\beta = \frac{N - Z}{N + Z}$$

- Taylor expansion

- Minimum at saturation density,  $\rho_0$
- Minimum in asymmetry:  $\beta = 0$
- Isospin symmetry  $\Rightarrow$  even powers of  $\beta$
- Give the coefficients a name!

$$\begin{aligned} \frac{E}{A}(\rho, \beta) &= \frac{E}{A}(\rho_0, \beta) \\ &\quad + 3\rho_0 \frac{\partial E/A}{\partial \rho} \Big|_{\rho_0} \left( \frac{\rho - \rho_0}{3\rho_0} \right) \\ &\quad + \frac{9\rho_0^2}{2!} \frac{\partial^2 E/A}{\partial \rho^2} \Big|_{\rho_0} \left( \frac{\rho - \rho_0}{3\rho_0} \right)^2 \\ &\quad + \mathcal{O}(3) \end{aligned}$$



# Effective approach to the EoS

- EoS provides a characterization of bulk properties:

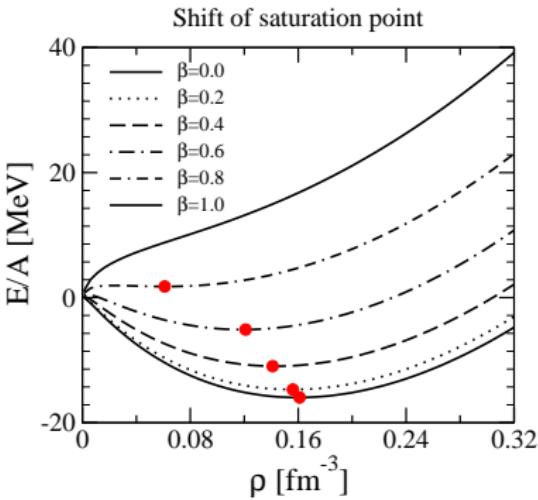
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- Minimum at saturation density,  $\rho_0$
- Minimum in asymmetry:  $\beta = 0$
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- Give the coefficients a name!

$$\begin{aligned}\frac{E}{A}(\rho, \beta) &= \frac{E}{A}(\rho_0, 0) + \frac{1}{2!} \left. \frac{\partial^2 E/A}{\partial \beta^2} \right|_{\rho_0, \beta=0} \beta^2 \\ &\quad + \frac{3\rho_0}{2!} \left. \frac{\partial^3 E/A}{\partial \beta^2 \partial \rho} \right|_{\rho_0, \beta=0} \beta^2 \left( \frac{\rho - \rho_0}{3\rho_0} \right) \\ &\quad + \frac{9\rho_0^2}{2!} \left\{ \left. \frac{\partial^2 E/A}{\partial \rho^2} \right|_{\rho_0, \beta=0} + \frac{1}{2!} \left. \frac{\partial^4 E/A}{\partial \rho^2 \beta^2} \right|_{\rho_0, \beta=0} \beta^2 \right\} \left( \frac{\rho - \rho_0}{3\rho_0} \right)^2 \\ &\quad + \mathcal{O}(3, 2)\end{aligned}$$



# Effective approach to the EoS

- EoS provides a characterization of bulk properties:

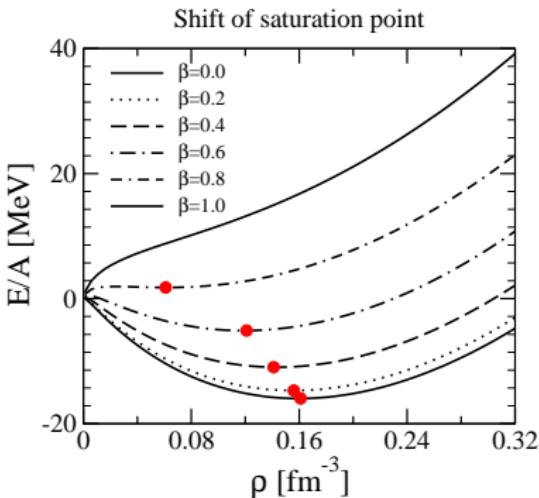
$$p(\varepsilon) = ? \iff \frac{E}{A}(\rho, \beta) = ?$$

$$\beta = \frac{N - Z}{N + Z}$$

- Taylor expansion

- Minimum at saturation density,  $\rho_0$
- Minimum in asymmetry:  $\beta = 0$
- Isospin symmetry  $\Rightarrow$  even powers of  $\beta$
- Give the coefficients a name!

$$\begin{aligned} \frac{E}{A}(\rho, \beta) &= E_0 + E_{sym} \beta^2 \\ &\quad + L \beta^2 \left( \frac{\rho - \rho_0}{3\rho_0} \right) \\ &\quad + \frac{1}{2!} \left\{ K_0 + K_{sym} \beta^2 \right\} \left( \frac{\rho - \rho_0}{3\rho_0} \right)^2 \\ &\quad + \mathcal{O}(3, 2) \end{aligned}$$



# EoS from basic nuclear properties

An incomplete list

$$\frac{E}{A}(\rho, \beta) = E_0 + E_{sym}\beta^2 + L\beta^2 \left( \frac{\rho - \rho_0}{3\rho_0} \right) + \frac{1}{2!} \{ K_0 + K_{sym}\beta^2 \} \left( \frac{\rho - \rho_0}{3\rho_0} \right)^2$$

Quantity	Experimental probes	Value	Ref.
$\rho_0$	$(e, e')$ elastic scattering	$0.16 \text{ fm}^{-3}$	[1]
$E_0$	$\frac{E}{A}$ bulk systematics	$-16 \text{ MeV}$	[1]
$K_0$	GMR energy in $Z \sim N$	$240 \pm 20 \text{ MeV}$	[2]
$E_{sym}$	$\frac{E}{A}$ bulk systematics + ID	$32 \pm 2 \text{ MeV}$	[3]
$L$	ID, IVMR energies, $\delta R$	$61 \pm 11 \text{ MeV}$	[3]
$K_{sym}$	?	?	

[1] Schuck & Ring, *The Nuclear Many-Body Problem* (Springer)

[2] Blaizot, Phys. Reps. 64, 171 (1981)

[3] Tsang *et al.*, Phys. Rev. Lett. 102, 122701 (2009)



## $^{208}\text{Pb}$ skin thickness & $L$

Neutron-matter pressure is dominated by  $L$ :

$$p(\rho_0, \beta) = \frac{\rho_0 \beta^2}{3} L$$

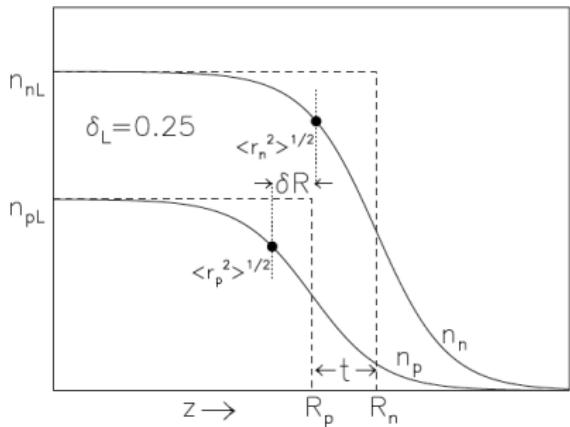
Does it correlate with nuclear observables?



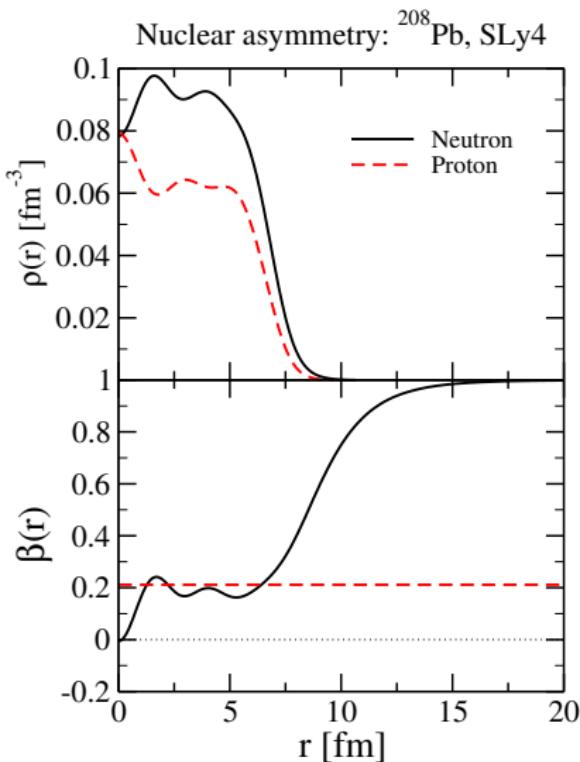
# Symmetry energy & nuclear structure

$$\text{<sup>208</sup>Pb skin thickness: } \delta R = R_n - R_p$$

## Neutron skin sketch



A. Steiner *et al.*, Phys. Rep. **411**, 325 (2005)

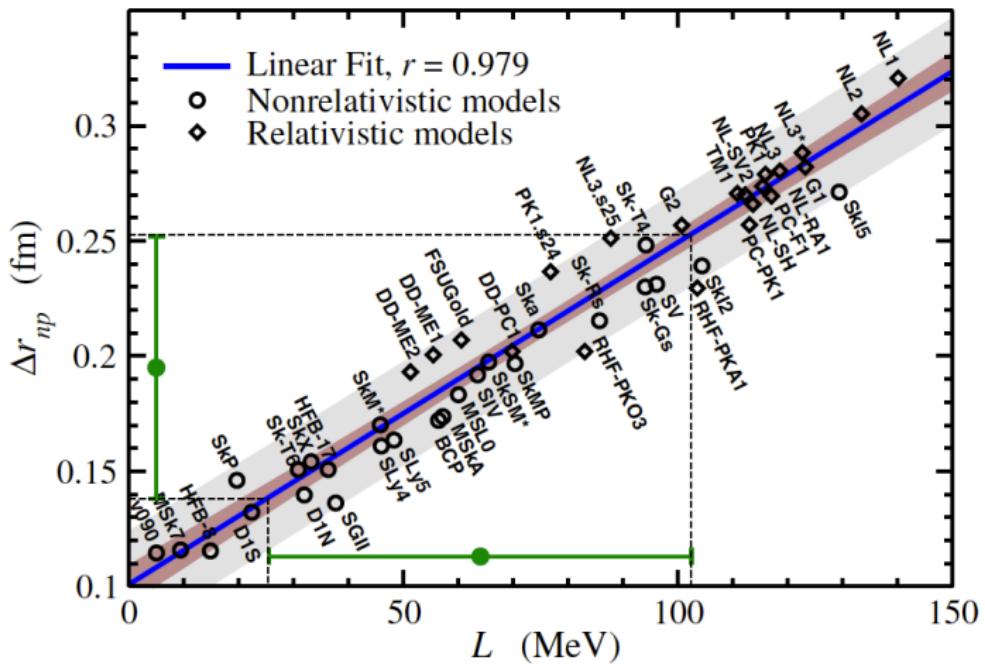


# Symmetry energy & nuclear structure



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### $\delta R$ vs $L$



S. Typel & B. Alex Brown, PRC 64, 027302 (2001)

X. Roca-Maza *et al.*, PRL 106, 252501 (2011)

# Symmetry energy & nuclear structure

## Parity violating electron scattering

$$q_n^W = -1$$

$$q_p^W = 1 - 4 \sin^2 \Theta_W \sim 0.05$$

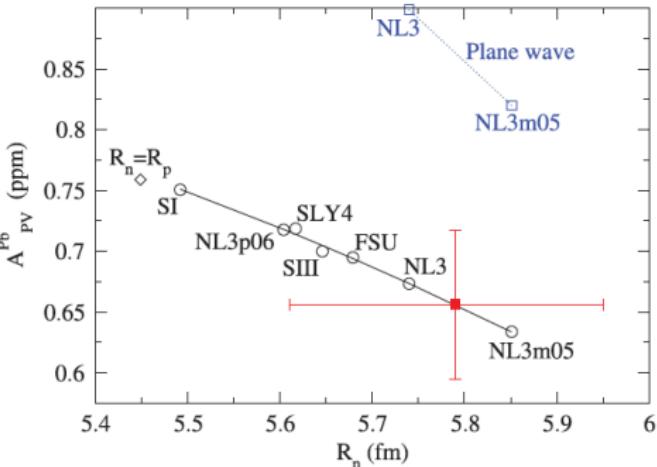
$$A_{\text{PM}} = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L}$$

Free of strong interaction uncertainties.

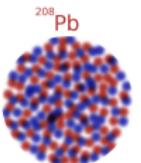
C. Horowitz *et al.*, PRC **63**, 025501 (2001)  
C. Horowitz *et al.*, PRC **85**, 032501(R) (2012)

## Lead Radius Experiment **PREX**

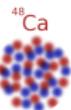
(early 2010 in Hall A)



Abrahamyan *et al.*, PRL **108**, 112502 (2012)



PREX



CREX

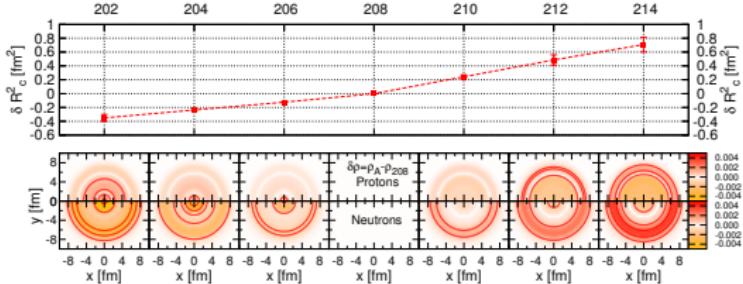
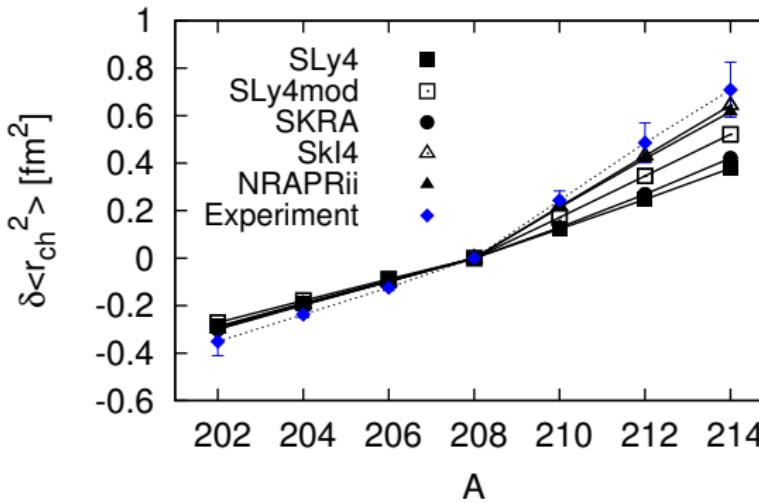
Edinburgh exp.:  $\pi^0$  photoproduction

# Kink in charge radii



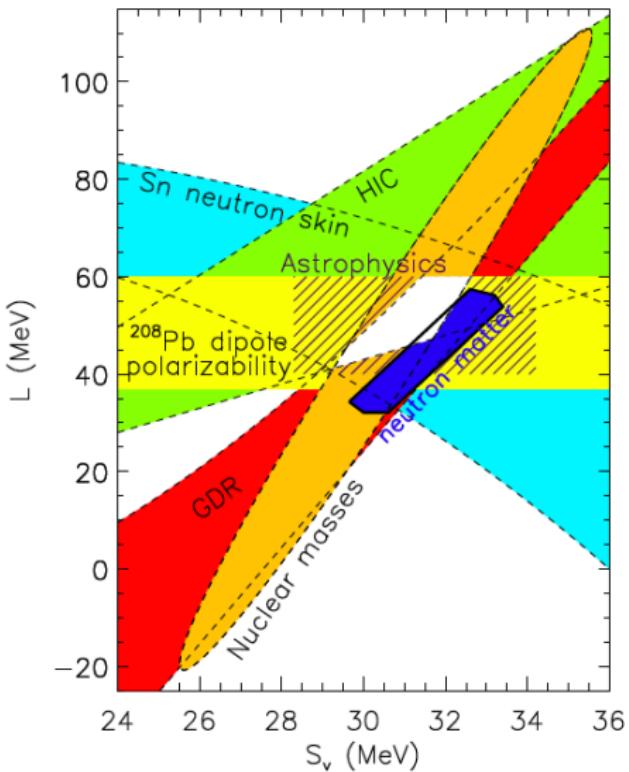
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## Radii of lead isotopes: open questions

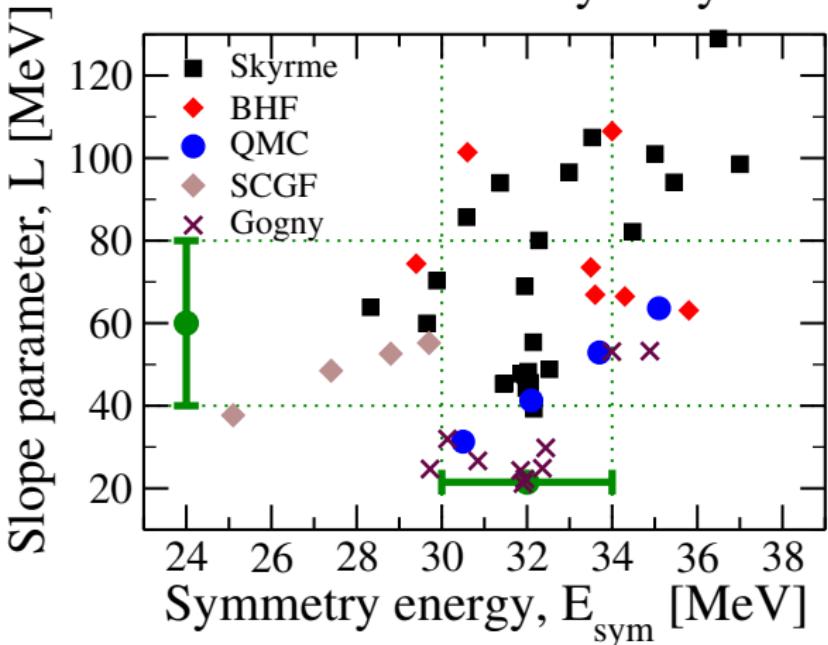


# Correlation between observables

## Correlations from different observables

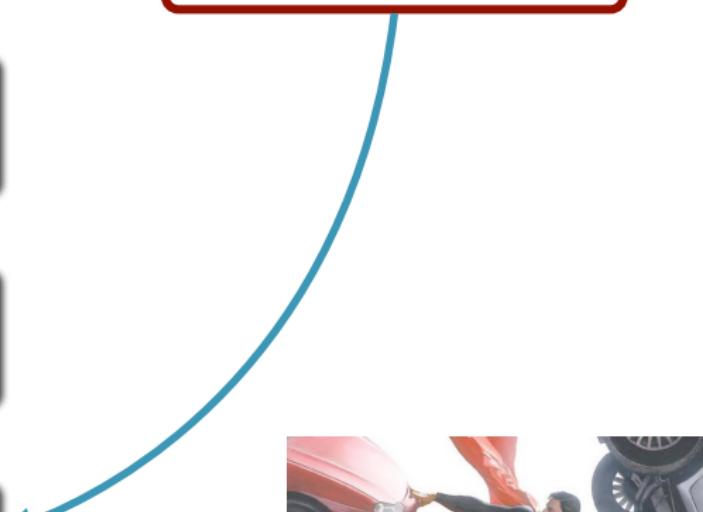
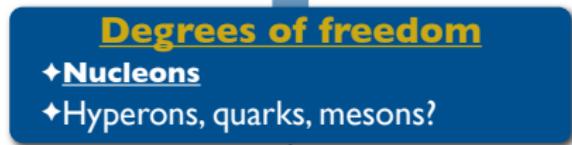


## Correlations from many-body theory



Gogny: Rosh Sellahewa, PhD



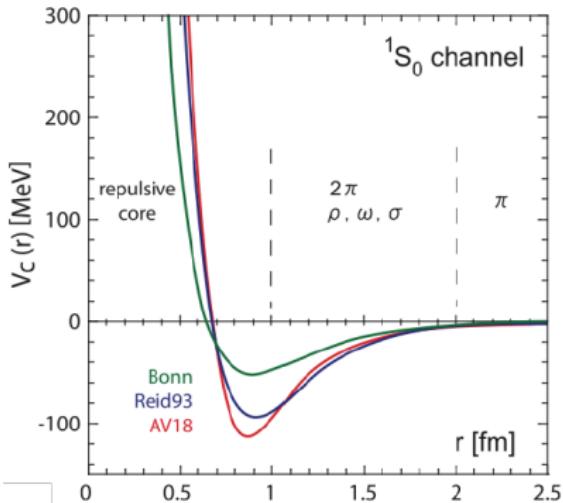


# Complications

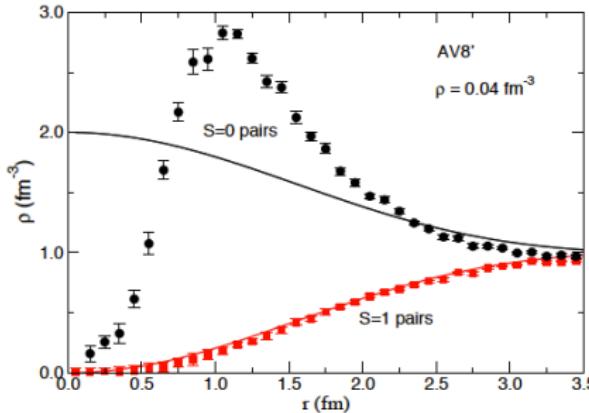
The hard life of nuclear many-body physicists

$$H = -\frac{\hbar^2}{2m} \sum_{i=1}^A \nabla_i^2 + \sum_{i < j} V_{ij} + \sum_{i < j < k} V_{ijk}$$

## Different NN potentials



## GFMC pair distribution function



Carlson *et al.*, Phys.Rev. C **68**, 025802 (2003)

- NN interaction is not uniquely defined
- Short-range core needs many-body treatment
- Chiral expansion  $\Rightarrow$  systematics & many-body forces



# Complications

The hard life of nuclear many-body physicists



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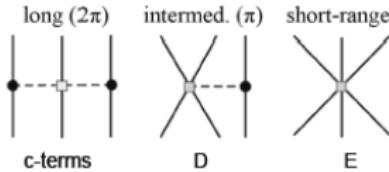
	NN	3N	4N
LO	X H	—	—
NLO	X H K X K H	—	—
N <sup>2</sup> LO	b H K	H X H X X	—
N <sup>3</sup> LO	X H K + ...	X H K + ...	X H + ...

Robert Roth – TU Darmstadt – 04/2013

## Chiral perturbation theory

- ①  $\pi$  and N as dof
- ② Systematic expansion in diagrams
- ③ 2N force at  $N^3LO$  - LEC are fitted
- ④ 3N force at  $N^2LO$  - 2 more LECs
- ⑤ 4N force are small

Tews, Schwenk et al., PRL 110 032504 (2013)



- NN interaction is not uniquely defined
- Short-range core needs many-body treatment
- Chiral expansion  $\Rightarrow$  systematics & many-body forces

## Diagrammatic expansion with 3BF

Self-energy expansion to 2nd order

$$\Sigma = \bullet \cdots \circ + \frac{1}{2} \text{ [diagram]} \quad \text{[diagram]}$$

The equation shows the definition of the self-energy operator  $\Sigma$ . It consists of a bare vertex (bullet) connected by a dashed line to a loop (circle). This is followed by a plus sign and half of a diagram where a bare vertex is connected to a loop via two dashed lines, with a vertical dashed line segment between them.

- Only skeleton 1PI diagrams needed

## Effective interaction expansion

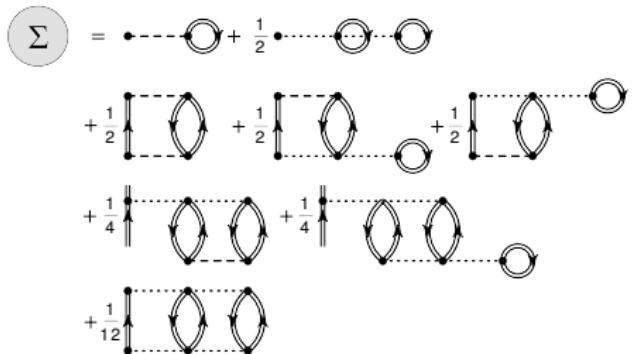
## Diagrammatic expansion with 3BF

### Self-energy expansion to 2nd order

$$\Sigma = \bullet \cdots \circlearrowleft + \frac{1}{2} \bullet \cdots \circlearrowleft \bullet \cdots \circlearrowright +$$

$$+ \frac{1}{2} \bullet \cdots \circlearrowleft \bullet \cdots \circlearrowright + \frac{1}{2} \bullet \cdots \circlearrowleft \bullet \cdots \circlearrowright + \frac{1}{2} \bullet \cdots \circlearrowleft \bullet \cdots \circlearrowright +$$

$$+ \frac{1}{4} \bullet \cdots \circlearrowleft \bullet \cdots \circlearrowright + \frac{1}{4} \bullet \cdots \circlearrowleft \bullet \cdots \circlearrowright +$$

$$+ \frac{1}{12} \bullet \cdots \circlearrowleft \bullet \cdots \circlearrowright$$


- Only skeleton 1PI diagrams needed
- Number of diagrams substantially increases
- Usable in higher order resummations

## Effective interaction expansion

### Define effective 1B and 2B forces

#### *Effective one-body force*

$$\bullet \cdots \times = \bullet \cdots \circlearrowleft + \frac{1}{2} \bullet \cdots \circlearrowright$$

#### *Effective two-body force*

$$\bullet \text{---} \bullet = \bullet \cdots \bullet + \bullet \cdots \bullet$$

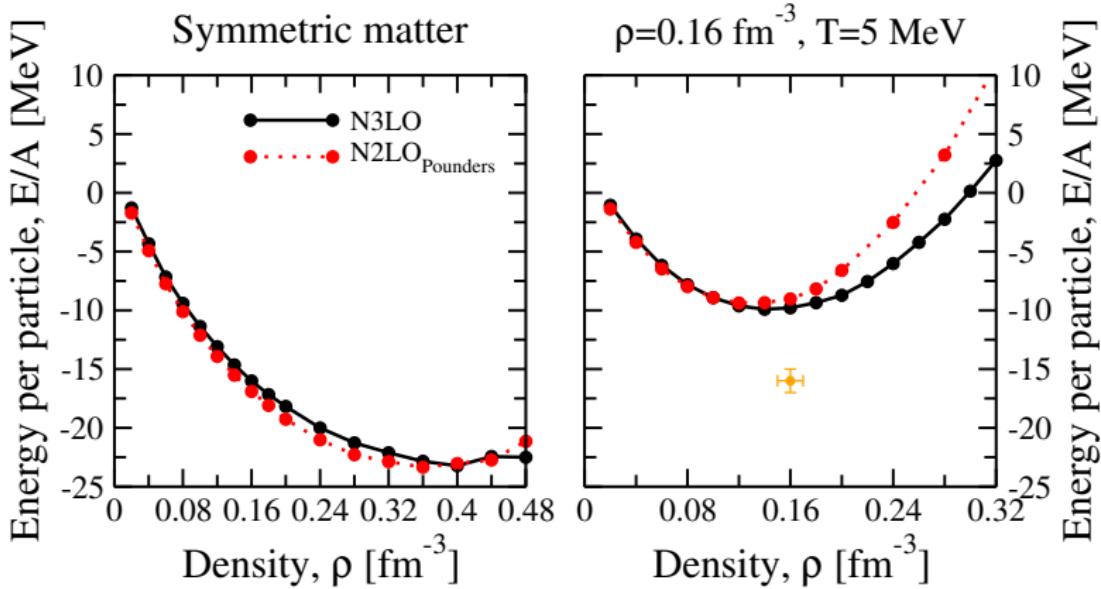
#### *Two-body propagator*

$$G_{\parallel} = \bullet \cdots \bullet + X + \bullet \text{---} \bullet$$

### Rewrite self-energy expansion

$$\Sigma = \bullet \cdots \times$$

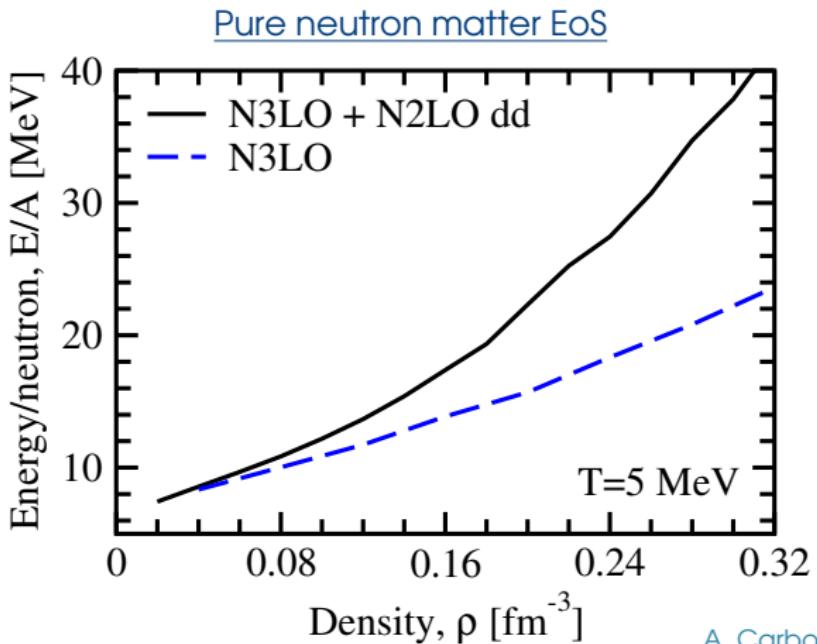
$$+ \frac{1}{2} \bullet \cdots \times + \frac{1}{12} \bullet \cdots \times$$



[A. Carbone](#), A. Rios & A. Polls

- Saturation properties of N3LO and NNLO similar
- NNLO  $\Rightarrow$  2B & 3B same order in  $\chi$  expansion
- Finite temperature & asymmetry also accessible

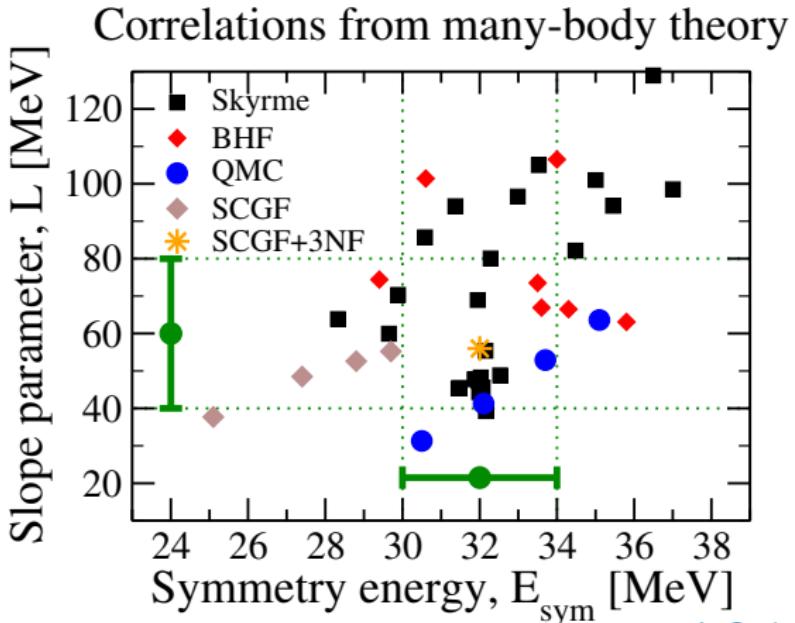




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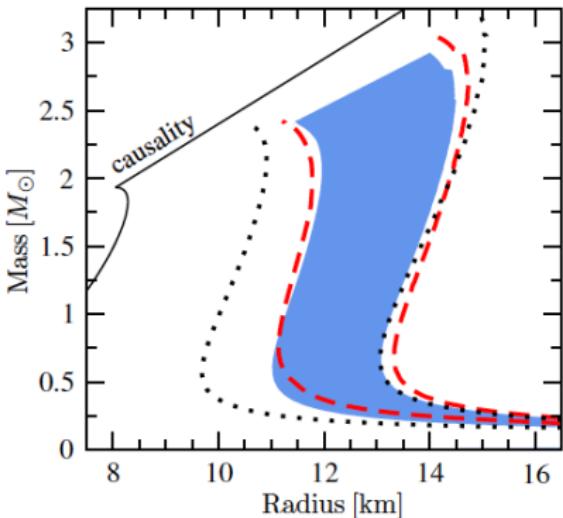
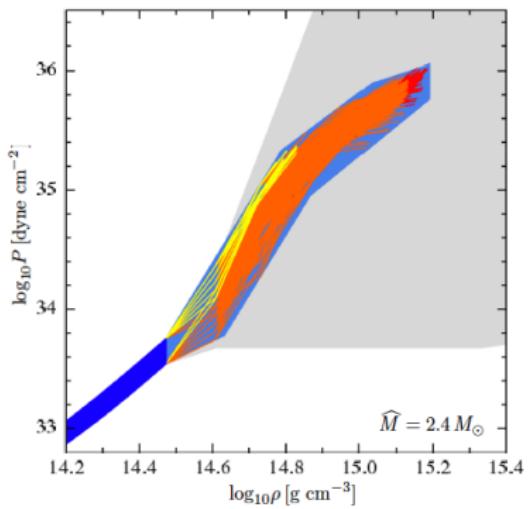


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## Impact on $M - R$ relation: causality + $2.4M_{\odot}$



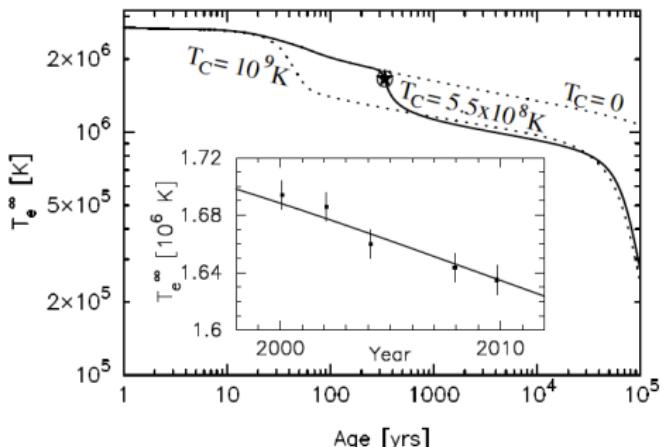
Hebeler, Lattimer, Pethick & Schwenk, arxiv:1303.4662

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# Pairing properties

Cooling of Cassiopea A &  $^3PF_2$  pairing



Page et al., PRL 106, 081101 (2011)

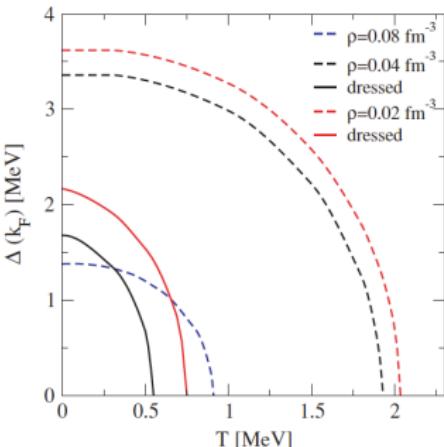
$$\Delta_{lk}^{JST} = - \sum_{l'} \int_0^\infty dk' k'^2 \langle kl | V^{JST} | k'l' \rangle \frac{\Delta_{l'k'}^{JST}}{2\chi_{k'}}$$

BCS / quasi-particle

$$\frac{1}{2\chi_k} = \frac{1 - 2f(\varepsilon_k)}{2\varepsilon_k}$$

$$\frac{1}{2\chi_k} = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \int_{-\infty}^{\infty} \frac{d\omega'}{2\pi} \frac{1 - f(\omega) - f(\omega')}{\omega + \omega'} A_k(\omega) A_k^s(\omega')$$

CDBonn:  $^1S_0$  pairing gap



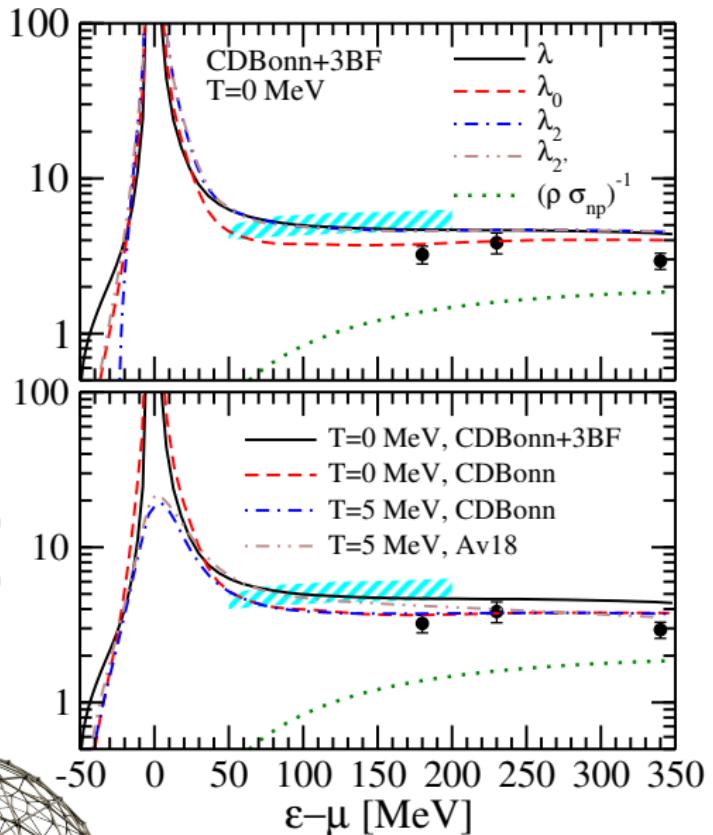
Muther & Dickhoff, PRC 72 054313 (2005)

Beyond quasi-particle



# Transport: nucleon mean-free path

Symmetric matter,  $\rho=0.16 \text{ fm}^{-3}$



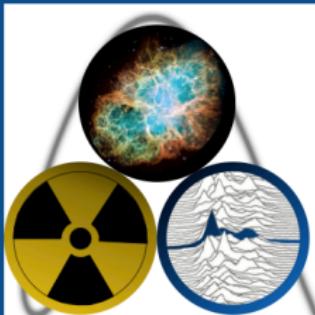
$$\lambda_k = \frac{1}{\Gamma_k} \frac{\partial \varepsilon_k}{\partial k}$$

- $\lambda \sim 4 - 5 \text{ fm}$  above 50 MeV
- Compatible with  $pA$  experiments
- Small model dependence
  - $\lambda_0 \Rightarrow$  no non-locality
  - $\lambda_2 \Rightarrow$  full non-locality
  - $\lambda_2' \Rightarrow m_k^* \text{ non-locality}$
- Classical approximation is incorrect!
- Little effect of 3BFs

- Ab initio description of nuclear & neutron matter
- Fully self-consistent & quantum mechanical calculation
- Single-particle microscopic properties ✓
- Thermodynamic properties ✓
- Mean-free path in dense matter ✓
- Adding three-body forces consistently
- Pairing:  $^1S_0$ ,  $^3PF_2$
- Other transport properties are coming



# Thanks!



## Neutron Stars

Nuclear Physics, Gravitational Waves & Astronomy

29-30 July 2013

Institute of Advanced Studies, University of Surrey

<http://www.ias.surrey.ac.uk/workshops/neutstar/>

V. Somà

T. U. Darmstadt



W. H. Dickhoff

Wash. U. St. Louis



A. Polls, A. Carbone

University of Barcelona



C. Barbieri, A. Cipollone

University of Surrey



Science & Technology  
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